

HOMOTOPY PERTURBATION METHOD TO SOLVE DUFFING-VAN DER POL EQUATION

BAGAYOGO Moussa^{1,*}, MINOUNGOU Youssouf², NEBIE Abdoul Wassiha³ and PARE Youssouf³

¹Centre Universitaire de Kaya Université Joseph KI-ZERBO Burkina Faso

²Ecole Normale Supérieure (ENS) Burkina Faso

³Université Joseph KI-ZERBO

Burkina Faso

Abstract

In this paper, the Homotopy Perturbation Method (HPM) and the Regular Perturbation Method (RPM) are used to study Duffing-Van der Pol equation. Then we compare the solutions obtained by these two methods.

Received: February 15, 2024; Accepted: April 13, 2024

2020 Mathematics Subject Classification: 74H10, 74H15, 35A35, 41A10, 65Yxx, 55PXX, 47F05.

Keywords and phrases: Duffing-Van der Pol equation, homotopy perturbation method, regular perturbation method.

*Corresponding author

How to cite this article: BAGAYOGO Moussa, MINOUNGOU Youssouf, NEBIE Abdoul Wassiha and PARE Youssouf, Homotopy perturbation method to solve Duffing-Van der Pol equation, Advances in Differential Equations and Control Processes 31(3) (2024), 299-315. https://doi.org/10.17654/0974324324016

This is an open access article under the CC BY license (http://creativecommons.org/licenses/by/4.0/).

Published Online: May 15, 2024

1. Introduction

The Duffing-Van der Pol equation which we examine is nonlinear and it contains two perturbation parameters μ and λ . The general form of this equation is:

$$\frac{d^2 u(t)}{dt^2} - \mu (1 - u^2(t)) \frac{du(t)}{dt} + u(t) + \lambda u^3(t) = f(t); \quad t > 0.$$
(1)

We here examine the following initial value problem:

$$\begin{cases} \frac{d^2 u(t)}{dt^2} - \mu(1 - u^2(t)) \frac{du(t)}{dt} + u(t) + \lambda u^3(t) = 0 \quad 0 < t < T \\ u(0) = 1; \qquad \qquad \frac{du(0)}{dt} = 0, \end{cases}$$
(2)

where $0 < \mu \ll 1$, $0 < \lambda \ll 1$.

First, we construct the solution of (2) by the HPM [4-6]. We then apply the regular perturbation method [2, 3] and finally we compare the two methods. We will be interested in the cases $\mu = \frac{1}{4}$ and $\mu = \frac{1}{8}$.

2. Application of the HPM to Solve the Duffing-Van der Pol Equation

In order to apply the HPM, we construct a homotopy H(v, p) which satisfies:

$$H(v, p) = (1 - p)[L(v) - L(u_0)] + p[L(v) + N(v) - f(t)] = 0$$

with

$$L(v) = \frac{d^2v}{dt^2} - \mu \frac{dv}{dt} + v, \ L(u_0) = \frac{d^2u_0}{dt^2} - \mu \frac{du_0}{dt} + u_0, \ N(v) = \mu v^2 \frac{dv}{dt} + \lambda v^3$$

and f(t) = 0.

Assuming

$$\frac{d^2 u_0}{dt^2} = \frac{du_0}{dt} = u_0 = 0,$$

we have

$$\frac{d^2v}{dt^2} - \mu \frac{dv}{dt} + v + p\mu v^2 \frac{dv}{dt} + p\lambda v^3 = 0.$$
(3)

Assume that the solution of equation (2) can be given by a power series in *p*:

$$v = v_0 + pv_1 + p^2 v_2 + \dots$$
 (4)

By substituting (4) into (3) and equating the terms with the identical powers of p, we have

$$p^{0}:\begin{cases} \frac{d^{2}v_{0}}{dt^{2}} - \mu \frac{dv_{0}}{dt} + v_{0} = 0\\ v_{0}(0) = 1; & \frac{dv_{0}(0)}{dt} = 0, \end{cases}$$
(5)

$$p^{1}:\begin{cases} \frac{d}{dt^{2}} - \mu \frac{dv_{1}}{dt} + v_{1} + \mu v_{0}^{2} \frac{dv_{0}}{dt} + \lambda v_{0}^{3} = 0\\ v_{1}(0) = 0; & \frac{dv_{1}(0)}{dt} = 0, \end{cases}$$
(6)

$$p^{2}:\begin{cases} \frac{d^{2}v_{2}}{dt^{2}} - \mu \frac{dv_{2}}{dt} + v_{2} + \mu v_{0}^{2} \frac{dv_{1}}{dt} + 2\mu v_{0} v_{1} \frac{dv_{0}}{dt} + 3\lambda v_{0}^{2} v_{1} = 0\\ v_{2}(0) = 0; \qquad \qquad \frac{dv_{2}(0)}{dt} = 0. \end{cases}$$
(7)

We solve the systems (5), (6), (7) for $\mu = \frac{1}{4}$ and $\mu = \frac{1}{8}$.

Case: $\mu = \frac{1}{4}$

From (5), we have

$$v_0(t) = e^{\frac{t}{8}} \left[\cos\left(\frac{3\sqrt{7}}{8}t\right) - \frac{1}{3\sqrt{7}}\sin\left(\frac{3\sqrt{7}}{8}t\right) \right].$$

From (6), we have

$$\begin{aligned} v_1(t) &= \left(-\frac{4355\lambda}{95634\sqrt{7}} - \frac{1843}{95634\sqrt{7}} \right) e^{\frac{3t}{8}} \sin\left(\frac{9\sqrt{7}t}{8}\right) \\ &+ \left(\frac{289\lambda}{10626} - \frac{109}{31878} \right) e^{\frac{3t}{8}} \cos\left(\frac{9\sqrt{7}t}{8}\right) \\ &+ \left[-\left(\frac{5}{126\sqrt{7}} + \frac{248\lambda}{63\sqrt{7}}\right) e^{\frac{3t}{8}} + \left(\frac{6541\lambda}{1518\sqrt{7}} + \frac{48}{253\sqrt{7}}\right) e^{\frac{t}{8}} \right] \sin\left(\frac{3\sqrt{7}t}{8}\right) \\ &+ \left[-\left(\frac{8\lambda}{21} + \frac{17}{126}\right) e^{\frac{3t}{8}} + \left(\frac{35}{253} + \frac{179\lambda}{506}\right) e^{\frac{t}{8}} \right] \cos\left(\frac{3\sqrt{7}t}{8}\right). \end{aligned}$$

From (7), we have

$$\begin{aligned} v_{2}(t) &= \left(-\frac{10638623\lambda^{2}}{4586989176\sqrt{7}} - \frac{1938703\lambda}{1720120941\sqrt{7}} + \frac{2718725}{13760967528\sqrt{7}} \right) e^{\frac{5t}{8}} \sin\left(\frac{15\sqrt{7}t}{8}\right) \\ &+ \left(\frac{38851\lambda^{2}}{72809352} - \frac{306961\lambda}{573373647} - \frac{350885}{4586989176} \right) e^{\frac{5t}{8}} \cos\left(\frac{15\sqrt{7}t}{8}\right) \\ &+ \left[\left(-\frac{17200306\lambda^{2}}{67278519\sqrt{7}} + \frac{28226119\lambda}{403671114\sqrt{7}} + \frac{223532807}{25834951296\sqrt{7}} \right) e^{\frac{5t}{8}} \right] \\ &+ \left(-\frac{5387645\lambda^{2}}{129042144\sqrt{7}} - \frac{235\lambda}{4554\sqrt{7}} + \frac{1666318497}{308\sqrt{7}} \right) e^{\frac{3t}{8}} \right] \sin\left(\frac{9\sqrt{7}t}{8}\right) \end{aligned}$$

$$\begin{split} &+ \left[\left(-\frac{12582\lambda^2}{118657} - \frac{1865485\lambda}{44852346} + \frac{1864883}{2870550144} \right) e^{\frac{5t}{8}} \right] \\ &+ \left(\frac{1000959\lambda^2}{14338016} + \frac{25\lambda}{506} + \frac{8278270321}{21252} \right) e^{\frac{3t}{8}} \right] \cos\left(\frac{9\sqrt{7}t}{8}\right) \\ &+ \left[\left(\frac{140948480\lambda^2}{67278519\sqrt{7}} + \frac{214983680\lambda}{201835557\sqrt{7}} + \frac{4426960}{201835557\sqrt{7}} \right) e^{\frac{5t}{8}} \right] \\ &- \left(\frac{1636\lambda^2}{483\sqrt{7}} + \frac{145\lambda}{66\sqrt{7}} + \frac{15803}{340032\sqrt{7}} \right) e^{\frac{3t}{8}} + \left(\frac{155643712503\lambda^2}{78361354016\sqrt{7}} \right) \\ &+ \frac{2788901\sqrt{7}\lambda}{19358042} - \frac{61289336254718701}{3716744064\sqrt{7}} \right) e^{\frac{t}{8}} \right] \sin\left(\frac{3\sqrt{7}t}{8}\right) \\ &+ \left[\left(-\frac{975872\lambda^2}{1067913} + \frac{6568960\lambda}{67278519} + \frac{1869232}{67278519} \right) e^{\frac{5t}{8}} \right] \\ &+ \left(\frac{4044\lambda^2}{1771} - \frac{25\lambda}{506} - \frac{17071}{340032} \right) e^{\frac{3t}{8}} + \left(\frac{104530270821\lambda^2}{78361354016} \right) \\ &+ \frac{1074589\lambda}{19358042} + \frac{482593173192685}{1238914688} \right) e^{\frac{t}{8}} \right] \cos\left(\frac{3\sqrt{7}t}{8}\right). \end{split}$$

Finally, the approximate solution of (2) is given by

$$u(t) \simeq \lim_{p \to 1} [v_0(t) + pv_1(t) + p^2 v_2(t)]$$
$$\simeq v_0(t) + v_1(t) + v_2(t).$$

Case: $\mu = \frac{1}{8}$

From (5), we have

$$v_0(t) = e^{\frac{t}{16}} \left[\cos\left(\frac{\sqrt{255}}{16}t\right) - \frac{1}{\sqrt{255}} \sin\left(\frac{\sqrt{255}}{16}t\right) \right].$$

From (6), we have

$$\begin{aligned} v_1(t) &= -\left(\frac{72707\lambda}{520710\sqrt{255}} + \frac{31939}{520710\sqrt{255}}\right) e^{\frac{3t}{16}} \sin\left(\frac{3\sqrt{255}t}{16}\right) \\ &+ \left(\frac{5249\lambda}{173570} - \frac{89}{104142}\right) e^{\frac{3t}{16}} \cos\left(\frac{3\sqrt{255}t}{16}\right) \\ &+ \left[-\left(\frac{4064\lambda}{85\sqrt{255}} + \frac{7\sqrt{3}}{170\sqrt{85}}\right) e^{\frac{3t}{16}} \right] \\ &+ \left(\frac{99901\lambda}{2042\sqrt{255}} + \frac{192\sqrt{3}}{1021\sqrt{85}}\right) e^{\frac{t}{16}} \right] \sin\left(\frac{\sqrt{255}t}{16}\right) \\ &+ \left[-\left(\frac{32\lambda}{85} + \frac{13}{102}\right) e^{\frac{3t}{16}} + \left(\frac{707\lambda}{2042} + \frac{131}{1021}\right) e^{\frac{t}{16}} \right] \cos\left(\frac{\sqrt{255}t}{16}\right). \end{aligned}$$

From (7), we have

$$v_{2}(t) = \left(-\frac{3141581903\lambda^{2}}{407018178600\sqrt{255}} - \frac{723890191\lambda}{152631816975\sqrt{255}} + \frac{37776433}{244210907160\sqrt{255}}\right)e^{\frac{5t}{16}}\sin\left(\frac{5\sqrt{255t}}{16}\right) \\ + \left(\frac{117014853\lambda^{2}}{135672726200} - \frac{21801073\lambda}{152631816975}\right)$$

$$\begin{aligned} &-\frac{342689}{14365347480} e^{\frac{5t}{16}} \cos\left(\frac{5\sqrt{255t}}{16}\right) \\ &+ \left[\left(-\frac{675884312\lambda^2}{163763295\sqrt{255}} + \frac{1449273611\lambda}{6878058390\sqrt{255}} \right. \\ &+ \frac{42452231179}{1760782947840\sqrt{255}} \right] e^{\frac{5t}{16}} \\ &+ \left(\frac{1523267\lambda^2}{347140\sqrt{255}} - \frac{25021\lambda}{173570\sqrt{255}} - \frac{2879\sqrt{3}}{347140\sqrt{85}} \right) e^{\frac{3t}{16}} \right] \sin\left(\frac{3\sqrt{255t}}{16}\right) \\ &+ \left[\left(-\frac{31203544\lambda^2}{272938825} - \frac{174224063\lambda}{3821143550} + \frac{4564663}{34525155840} \right) e^{\frac{5t}{16}} \right] \\ &+ \left[\left(-\frac{32829\lambda^2}{347140} + \frac{8253\lambda}{173570} + \frac{61}{347140} \right) e^{\frac{3t}{16}} \right] \cos\left(\frac{3\sqrt{255t}}{16}\right) \\ &+ \left[\left(\frac{4251910144\lambda^2}{163763295\sqrt{255}} + \frac{41970122752\lambda}{3439029195\sqrt{255}} \right) \\ &+ \frac{227918144}{3439029195\sqrt{255}} \right] e^{\frac{5t}{16}} \\ &- \left(\frac{1172528\lambda^2}{86785\sqrt{85}} + \frac{1422529\sqrt{3}\lambda}{173570\sqrt{85}} + \frac{2987201}{22216960\sqrt{255}} \right) e^{\frac{3t}{16}} \\ &+ \left(\frac{3051545301\sqrt{3}\lambda^2}{694794584\sqrt{85}} + \frac{4877576339\sqrt{3}\lambda}{1215890522\sqrt{85}} \right) \\ &+ \frac{20650444499}{311267973632\sqrt{255}} \right] e^{\frac{t}{16}} \\ \end{bmatrix}$$

$$+ \left[\left(-\frac{1169752064\lambda^2}{272938825} + \frac{1528152064\lambda}{17195145975} + \frac{4954816}{202295835} \right) e^{\frac{5t}{16}} \right] \\ + \left(\frac{783504\lambda^2}{86785} - \frac{8253\lambda}{173570} - \frac{1059519}{22216960} \right) e^{\frac{3t}{16}} \\ + \left(\frac{3281824383\lambda^2}{694794584} + \frac{52445677\lambda}{1215890522} \right) \\ - \frac{1018853333}{44466853376} e^{\frac{t}{16}} \right] \cos \left(\frac{\sqrt{255t}}{16} \right).$$

Finally, the approximate solution of (2) is given by

$$u(t) \simeq \lim_{p \to 1} [v_0(t) + pv_1(t) + p^2 v_2(t)]$$
$$\simeq v_0(t) + v_1(t) + v_2(t).$$

3. The Regular Perturbation Method

Let us suppose that the solution u(t) of the initial value problem (2) has the following form [2]:

$$u(t) = \sum_{n=0}^{+\infty} \lambda^n u_n(t) + R_N(t, \lambda), \qquad (8)$$

where $R_N(t, \lambda)$ is the remainder of the series.

Taking (8) into (2), and collecting equal powers of λ we obtain a system of recurrent initial value problems for $u_n(t)$, n = 0, 1, 2, ...

$$\lambda^{0} : \begin{cases} \frac{d^{2}u_{0}}{dt^{2}} + u_{0} - \mu(1 - u_{0}^{2})\frac{du_{0}}{dt} = 0\\ u_{0}(0) = 1; & \frac{du_{0}(0)}{dt} = 0, \end{cases}$$
(9)

Solving of Duffing-Van der Pol Equation

$$\lambda^{1} : \begin{cases} \frac{d^{2}u_{1}}{dt^{2}} + u_{1} - \mu \left[(1 - u_{0}^{2}) \frac{du_{1}}{dt} - 2u_{0}u_{1} \frac{du_{0}}{dt} \right] + u_{0}^{3} = 0 \\ u_{1}(0) = 0; & \frac{du_{1}(0)}{dt} = 0, \end{cases}$$
(10)
$$\lambda^{2} : \begin{cases} \frac{d^{2}u_{2}}{dt^{2}} + u_{2} - \mu \left[(1 - u_{0}^{2}) \frac{du_{2}}{dt} - 2u_{0}u_{1} \frac{du_{1}}{dt} \\ - (u_{1}^{2} + 2u_{0}u_{2}) \frac{du_{0}}{dt} \right] + 3u_{0}^{2}u_{1} = 0 \\ u_{2}(0) = 0; & \frac{du_{2}(0)}{dt} = 0. \end{cases}$$
(11)

Consider the initial value problems (9), (10) and (11), and suppose that the solution of each one of them is of the following form:

$$u_k(t) = \sum_{n=0}^{+\infty} \mu^n u_{kn}(t),$$
(12)

where k = 0, 1, 2.

Taking (12) into (9), (10) and (11), and collecting equal powers of μ , we obtain a linear system of recurrent initial value problems for $u_{kn}(t)$, k = 0, 1, 2 and n = 0, 1, ...

For k = 0

$$\mu^{0} : \begin{cases} \frac{d^{2}u_{00}}{dt^{2}} + u_{00} = 0 \\ u_{00}(0) = 1; & \frac{du_{00}(0)}{dt} = 0, \end{cases}$$
(13)
$$\mu^{1} : \begin{cases} \frac{d^{2}u_{01}}{dt^{2}} + u_{01} - (1 - u_{00}^{2})\frac{du_{00}}{dt} = 0 \\ u_{01}(0) = 1; & \frac{du_{01}(0)}{dt} = 0. \end{cases}$$
(14)

For k = 1

$$\mu^{0} : \begin{cases} \frac{d^{2}u_{10}}{dt^{2}} + u_{10} + u_{00}^{3} = 0 \\ u_{10}(0) = 0; & \frac{du_{10}(0)}{dt} = 0, \end{cases}$$
(15)
$$\mu^{1} : \begin{cases} \frac{d^{2}u_{10}}{dt^{2}} + u_{11} - \left[(1 - u_{00}^{2}) \frac{du_{10}(0)}{dt} \\ - 2u_{00}u_{10} \frac{du_{00}(0)}{dt} - 3u_{00}^{2}u_{01} \right] = 0 \\ u_{11}(0) = 0; & \frac{du_{11}(0)}{dt} = 0. \end{cases}$$
(16)

For k = 2

$$\mu^{0} : \begin{cases} \frac{d^{2}u_{20}}{dt^{2}} + u_{20} + 3u_{00}^{3}u_{10} = 0 \\ u_{20}(0) = 0; & \frac{du_{20}(0)}{dt} = 0, \end{cases}$$
(17)
$$\mu^{1} : \begin{cases} \frac{d^{2}u_{21}}{dt^{2}} + u_{21} - \left[(1 - u_{00}^{2})\frac{du_{20}}{dt} - 2u_{00}u_{10}\frac{du_{10}}{dt} \\ - (u_{10}^{2} + 2u_{00}u_{20})\frac{du_{00}}{dt} + 3u_{00}(2u_{01}u_{10} + u_{00}u_{11}) \right] = 0 \\ u_{20}(0) = 0; & \frac{du_{20}(0)}{dt} = 0. \end{cases}$$

From (13), we have

$$u_{00}(t) = \cos t.$$

From (14), we have

$$u_{01}(t) = -\frac{1}{32}\sin(3t) - \frac{9}{32}\sin t + \frac{3}{8}t\cos t.$$

From (15), we have

$$u_{10}(t) = \frac{1}{32}\cos(3t) - \frac{3}{8}t\sin t - \frac{1}{32}\cos t.$$

From (16), we have

$$u_{11}(t) = -\frac{1}{384}\sin(5t) - \frac{9}{256}\sin(3t) + \left(-\frac{9}{32}t^2 + \frac{295}{768}\right)\sin t - \frac{17}{64}t\cos t.$$

From (17), we have

$$u_{20}(t) = \frac{1}{1024}\cos(5t) - \frac{9}{256}t\sin(3t) - \frac{3}{128}\cos(3t) + \frac{3}{32}t\sin t + \left(-\frac{9}{128}t^2 + \frac{23}{1024}\right)\cos t.$$

From (18), we have

$$u_{21}(t) = -\frac{7}{49152}\sin(7t) + \frac{11}{49152}\sin(5t) - \frac{25}{8192}t\cos(5t) + \left(-\frac{135}{4096}t^2 + \frac{361}{4096}\right)\sin(3t) - \frac{717}{8192}t\cos(3t) + \left(\frac{915}{4096}t^2 - \frac{4209}{8192}\right)\sin t + \left(-\frac{81}{1024}t^3 + \frac{87}{256}t\right)\cos t.$$

We give the approximate solution of (2) for $\mu = \frac{1}{4}$ and $\mu = \frac{1}{8}$.

Case: $\mu = \frac{1}{4}$ $u(t) \approx u_0(t) + \lambda u_1(t) + \lambda^2 u_2(t)$ $\approx \left[u_{00}(t) + \frac{1}{4} u_{01}(t) \right] + \left[u_{10}(t) + \frac{1}{4} u_{11}(t) \right] \lambda + \left[u_{20}(t) + \frac{1}{4} u_{21}(t) \right] \lambda^2.$

Case:
$$\mu = \frac{1}{8}$$

 $u(t) \simeq u_0(t) + \lambda u_1(t) + \lambda^2 u_2(t)$
 $\simeq \left[u_{00}(t) + \frac{1}{8} u_{01}(t) \right] + \left[u_{10}(t) + \frac{1}{8} u_{11}(t) \right] \lambda + \left[u_{20}(t) + \frac{1}{8} u_{21}(t) \right] \lambda^2$

4. Solutions Analysis

In this section, we analyze the approximate solutions of (2) obtained by the two numerical methods (HPM and RPM). We vary the perturbation parameters μ and λ in order to observe their incidence on the approximate solutions.

Tables 1 and 2 give some values of the solutions approached for $\mu = \frac{1}{4}$, $\mu = \frac{1}{8}$, $\mu = \frac{1}{16}$ and $\mu = \frac{1}{32}$.

Table 1. Numerical analysis for $\lambda = 0.002$

(a) $\mu = \frac{1}{4}$				
t	U _{hpm}	U _{rpm}	$\left U_{hpm} - U_{rpm} \right $	
0	1.0000	1.0000	0	
0.1	0.9950	0.9950	0	
0.2	0.9800	0.9800	0	
0.3	0.9552	0.9552	0	
0.4	0.9208	0.9208	0	
0.5	0.8770	0.8770	0	
0.6	0.8241	0.8241	0	
0.7	0.7625	0.7626	0.0001	
0.8	0.6927	0.6927	0	
0.9	0.6149	0.6150	0.0001	
1	0.5297	0.5299	0.0002	

		8	
t	U_{hpm}	U _{rpm}	$U_{hpm} - U_{rpm}$
0	1.0000	1.0000	0
0.1	0.9950	0.9950	0
0.2	0.9800	0.9800	0
0.3	0.9552	0.9552	0
0.4	0.9208	0.9208	0
0.5	0.8772	0.8772	0
0.6	0.8246	0.8246	0
0.7	0.7635	0.7635	0
0.8	0.6944	0.6944	0
0.9	0.6180	0.6180	0
1	0.5347	0.5348	0.0001

(b) $\mu = \frac{1}{8}$

Table 2. Numerical	analysis for	$\lambda = 0.002$
--------------------	--------------	-------------------

		10	
t	U_{hpm}	U _{rpm}	$U_{hpm} - U_{rpm}$
0	1.0000	1.0000	0
0.1	0.9950	0.9950	0
0.2	0.9800	0.9800	0
0.3	0.9552	0.9552	0
0.4	0.9209	0.9209	0
0.5	0.8773	0.8773	0
0.6	0.8248	0.8248	0
0.7	0.7640	0.7640	0
0.8	0.6953	0.6953	0
0.9	0.6195	0.6195	0
1	0.5371	0.5372	0.0001

(a)
$$\mu = \frac{1}{16}$$

	32		
t	U_{hpm}	U _{rpm}	$U_{hpm} - U_{rpm}$
0	1.0000	1.0000	0
0.1	0.9950	0.9950	0
0.2	0.9800	0.9800	0
0.3	0.9552	0.9552	0
0.4	0.9209	0.9209	0
0.5	0.8773	0.8773	0
0.6	0.8249	0.8249	0
0.7	0.7642	0.7642	0
0.8	0.6958	0.6958	0
0.9	0.6202	0.6202	0
1	0.5384	0.5384	0

(b) $\mu = \frac{1}{32}$

It is noted that the variation of the parameter λ does not influence the approximate solutions so much. However the variation of the parameter μ disturbs the approximate solutions. For small time interval the solutions obtained by the two methods are practically the same. But for large time interval the solutions obtained by the two methods are different. We can observe it on Figures 2-4.



Figure 1. Comparison of the HPM solution with RPM solution for $\mu = \frac{1}{4}$.



Figure 2. Comparison of the HPM solution with RPM solution for $\mu = \frac{1}{8}$.



Figure 3. Comparison of the HPM solution with RPM solution for $\mu = \frac{1}{16}$.



Figure 4. Comparison of the HPM solution with RPM solution for $\mu = \frac{1}{16}$.



Figure 5. Solution of the HPM and RPM solution for different values of λ .

5. Conclusion

In this paper, we applied HPM to obtain analytical solution of Duffing-Van der Pol equation. The results obtained from this method have been compared with those obtained from regular perturbation. A numerical comparison between HPM and regular perturbation method is depicted in Figures 2-4. The HPM converges in a small time interval and computationally takes long time for a large interval. The HPM series gives reasonable results in the small time interval.

References

- [1] A. Beléndez, A. Hernández, T. Beléndez, E. Fernández, M. L. Álvarez and C. Neipp, Application of He's homotopy perturbation method to the Duffing-Harmonic oscillator, International Journal of Nonlinear Sciences and Numerical Simulation 8(1) (2007), 79-88.
- [2] E. M. de Jager and Jiang Furu, The Theory of Singular Perturbations, North-Holland, 1996.
- [3] Gabriel BISSANGA, Application of Adomian decomposition method to solving the Duffing equation, Comparison with perturbation method, Proceedings of the

Fourth International Workshop on Contemporary Problems in Mathematical Physics, Cotonou Benin, November 2005. World Scientific Publishing Co. Pte. Ltd., 2006.

- [4] J. H. He, Homotopy perturbation technique, Computer Methods in Applied Mechanics and Engineering 178 (1999), 257-262.
- [5] J. H. He, Homotopy perturbation method: a new nonlinear analytical technique, Applied Mathematics and Computation 135 (2003), 73-79.
- [6] J.-H. He, The homotopy perturbation method for nonlinear oscillators with discontinuities, Applied Mathematics and Computation 151(1) (2004), 287-292.
- [7] A. Kimiaeifar, A. R. Saidi, G. H. Bagheri and M. Rahimpour, Analytical solution for Van der Pol-Duffing oscillators, Chaos Solitons Fractals 42(5) (2009), 2660-2666.
- [8] Rodrigo D. Euzébio and Jaume Llibre, Sufficient conditions for the existence of periodic solutions of the extended Duffing-Van der Pol oscillator, International Journal of Computer Mathematics 93(8) (2016), 1358-1382.
- [9] H. Mirgolbabaee, S. T. Ledari and D. D. Gaji, New approach method for solving Duffing-type nonlinear oscillator, Alexandria Engineering Journal 55(2) (2016), 1695-1702.
- [10] M. Momeni, N. Jamshidi, A. Barari and D. D. Ganji, Application of He's energy balance method to Duffing harmonic oscillators, International Journal of Computer Mathematics 88(1) (2011), 135-144.
- [11] Aiyong Chen and Guirong Jiang, Periodic solution of the Duffing-Van der Pol oscillator by homotopy perturbation method, International Journal of Computer Mathematics 87(12) (2010), 2688-2696.
- [12] J. Kyzioł and A. Okniński, The Duffing-Van der Pol equation: Metamorphoses of resonance curves, Nonlinear Dynamics and Systems Theory 15 (2015), 25-31.