

DEVELOPING ANSATZ METHOD FOR SOLVING THE NEUTRON DIFFUSION SYSTEM UNDER GENERAL PHYSICAL CONDITIONS

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Abstract

This paper develops an ansatz method to derive the analytic solution of the neutron flux system under general initial conditions. Explicit closed series forms are established for the neutron flux and the delayed neutron concentration in terms of exponential functions. Existing results in the literature are recovered as special cases.

1. Introduction

In this paper, we consider the following system:

$$\frac{1}{V}\frac{\partial\phi}{\partial t} = D\frac{\partial^2\phi}{\partial x^2} + \left(-\sum_a + (1-\beta)\nu\sum_f\right)\phi(x,t) + \lambda C(x,t), \quad (1)$$

$$\frac{\partial C}{\partial t} = \beta v \sum_{f} \phi(x, t) - \lambda C(x, t), \qquad (2)$$

where $\phi(x, t)$ and C(x, t) stand for the neutron flux and the delayed neutron concentration, respectively. The parameters V, \sum_{a} , β , ν , \sum_{f} and λ in the system (1)-(2) were well defined/described in [1, 2]. The system is governed by the boundary conditions (BCs):

$$\phi(0, t) = 0, \quad \phi(L, t) = 0, \quad t > 0, \tag{3}$$

and the following initial conditions (ICs) in general form

$$\phi(x, 0) = f(x), \quad C(x, 0) = g(x), \quad 0 < x < L, \tag{4}$$

where f(x) and g(x) are given real functions. The model (1)-(4) is of great importance in both reactor design and theoretical nuclear physics. Obtaining accurate solution for the above system is essential for the purpose of safety considerations. In the literature, many authors [1-8] employed different techniques to solve the present system under specific constant ICs. These authors implemented f(x) and g(x) as constant functions, given by [1-8]:

$$f(x) = \phi_0, \quad g(x) = \frac{\beta \nu \sum_f}{\lambda} \phi_0, \quad 0 < x < L.$$
 (5)

Although the authors [1-8] introduced different effective methods to solve the neutron diffusion system, such methods ignored the sense to obtain the solution in a direct manner. The preceding discussion formed the main purpose of the present work. Very recently, Al-Sharif et al. [9] developed ansatz method to solve the system (1)-(2) under the BCs (3) and the special case of the ICs (5). In [9], the authors showed that their ansatz method enjoyed many advantages over the previous analytical and numerical approaches [1-8]. As far as we know, the literature is rich of various methods to solve ordinary differential equations (ODEs) and partial differential equations (PDEs) such as the LT [10-18], differential transform method (DTM) [19, 20], the homotopy analysis method (HAM) [21, 22], the homotopy perturbation method (HPM) [23-27], and the Adomian decomposition method (ADM) [28-31]. However, most of the above methods need massive calculations to reach the desired accurate solution. The objective of this work is to introduce a simple procedure to obtain the solution of the system (1)-(2) under the BCs (3) and the general ICs (4).

2. Analytic Approach

Let us rewrite equations (1) and (2) as

$$\frac{\partial \phi}{\partial t} = VD \frac{\partial^2 \phi}{\partial x^2} + \omega \phi(x, t) + \lambda VC(x, t), \tag{6}$$

$$\frac{\partial C}{\partial t} = \alpha \phi(x, t) - \lambda C(x, t), \tag{7}$$

where

$$\omega = V \left(-\sum_{a} + (1 - \beta) \nu \sum_{f} \right), \quad \alpha = \beta \nu \sum_{f}.$$
 (8)

Following Al-Sharif et al. [9], we assume $\phi(x, t)$ and C(x, t) in the forms:

$$\phi(x, t) = \sum_{n=0}^{\infty} \sin(\gamma_n x) (A_n e^{\varepsilon_n t} + B_n e^{\delta_n t}), \qquad (9)$$

$$C(x, t) = \sum_{n=0}^{\infty} \sin(\gamma_n x) \left(E_n e^{\varepsilon_n t} + F_n e^{\delta_n t} \right), \tag{10}$$

where $\gamma_n = (2n+1)\frac{\pi}{L}$. The coefficients A_n , B_n , E_n , ε_n , and δ_n are unknowns, to be evaluated later. It may be important here to mention that the ansatz/assumption (9) automatically satisfies the BCs (3). Substituting equations (9) and (10) into equations (6) and (7) yields the following system (see [9] for details):

$$(\varepsilon_n - \omega + VD\gamma_n^2)A_n - \lambda VE_n = 0, \tag{11}$$

$$(\delta_n - \omega + VD\gamma_n^2)B_n - \lambda VF_n = 0, \qquad (12)$$

$$(\varepsilon_n + \lambda)E_n - \alpha A_n = 0, \tag{13}$$

$$(\delta_n + \lambda)F_n - \alpha B_n = 0. \tag{14}$$

Applying the first IC in (3) implies

$$\sum_{n=0}^{\infty} \sin(\gamma_n x) (A_n + B_n) = f(x).$$
(15)

Thus as in [18],

$$A_n + B_n = \frac{2}{L} \int_0^L f(x) \sin(\gamma_n x) dx,$$
 (16)

or

$$A_n + B_n = I_f, (17)$$

where I_f is defined by

$$I_f = \frac{2}{L} \int_0^L f(x) \sin(\gamma_n x) dx.$$
(18)

Similarly, the second IC in (3) leads to

$$\sum_{n=0}^{\infty} \sin(\gamma_n x) (E_n + F_n) = g(x), \tag{19}$$

and, hence on account of [18],

$$E_n + F_n = I_g, (20)$$

where I_g is defined by

$$I_g = \frac{2}{L} \int_0^L g(x) \sin(\gamma_n x) dx.$$
 (21)

Accordingly, we have the algebraic system:

$$(\varepsilon_n - \omega + VD\gamma_n^2)A_n - \lambda VE_n = 0, \qquad (22)$$

$$(\delta_n - \omega + VD\gamma_n^2)B_n - \lambda VF_n = 0, \qquad (23)$$

$$(\varepsilon_n + \lambda)E_n - \alpha A_n = 0, \qquad (24)$$

$$(\delta_n + \lambda)F_n - \alpha B_n = 0, \qquad (25)$$

$$A_n + B_n = I_f, (26)$$

$$E_n + F_n = I_g. (27)$$

This system must be solved in order to determine the unknowns *A*, *B*, *E*, *F*, ε_n and δ_n . In [9], the authors showed that ε_n and δ_n can be obtained as

$$\varepsilon_n = \frac{a_n + \sqrt{a_n^2 + 4b_n}}{2}, \quad \delta_n = \frac{a_n - \sqrt{a_n^2 + 4b_n}}{2},$$
(28)

where

$$a_n = \omega - \lambda - VD\gamma_n^2, \quad b_n = \lambda(\omega + \alpha V - VD\gamma_n^2).$$
 (29)

From (22) and (23), we have

$$A_n = \frac{\lambda V E_n}{\varepsilon_n - \Omega_n}, \quad B_n = \frac{\lambda V F_n}{\delta_n - \Omega_n}, \tag{30}$$

where

$$\Omega_n = \omega - V D \gamma_n^2. \tag{31}$$

Substituting (30) into (26), we get

$$\frac{\lambda V E_n}{\varepsilon_n - \Omega_n} + \frac{\lambda V F_n}{\delta_n - \Omega_n} = I_f.$$
(32)

On solving (32) and (27) for E_n and simplifying, we obtain

$$E_n = \frac{(\varepsilon_n - \Omega_n)[(\delta_n - \Omega_n)I_f - \lambda VI_g]}{\lambda V(\delta_n - \varepsilon_n)}, \quad \delta_n \neq \varepsilon_n.$$
(33)

We can prove that $(\varepsilon_n - \Omega_n)(\delta_n - \Omega_n) = -\alpha\lambda V$ (see Appendix in [9]). Therefore, equation (33) becomes

$$E_n = \frac{\alpha I_f + (\varepsilon_n - \Omega_n) I_g}{\varepsilon_n - \delta_n}, \quad \delta_n \neq \varepsilon_n.$$
(34)

Employing (34) in (27),

$$F_n = \frac{-\alpha I_f + (\Omega_n - \delta_n) I_g}{\varepsilon_n - \delta_n}, \quad \delta_n \neq \varepsilon_n.$$
(35)

From (34) and (35), it is obvious that the sum $E_n + F_n$ equals I_g which is equivalent to equation (27). Inserting (34) and (35) into equation (30), we obtain A_n and B_n as

$$A_n = \frac{\lambda V[\alpha I_f + (\varepsilon_n - \Omega_n) I_g]}{(\varepsilon_n - \Omega_n)(\varepsilon_n - \delta_n)}$$
(36)

and

$$B_n = \frac{\lambda V[(\Omega_n - \delta_n)I_g - \alpha I_f]}{(\delta_n - \Omega_n)(\varepsilon_n - \delta_n)},$$
(37)

respectively. From the equality $(\varepsilon_n - \Omega_n)(\delta_n - \Omega_n) = -\alpha\lambda V$, we have $\varepsilon_n - \Omega_n = -\alpha\lambda V/(\delta_n - \Omega_n)$ and hence can get A_n in the form:

$$A_n = \frac{\lambda V I_g - (\delta_n - \Omega_n) I_f}{\varepsilon_n - \delta_n}.$$
(38)

Similarly, B_n takes the form

$$B_n = \frac{-\lambda V I_g + (\varepsilon_n - \Omega_n) I_f}{\varepsilon_n - \delta_n}.$$
(39)

It may be noted from equations (38) and (39) that $A_n + B_n = I_f$ which is equivalent to equation (26).

2.1. The analytic solution

From equation (9), we obtain $\phi(x, t)$ as

$$\phi(x, t) = \sum_{n=0}^{\infty} \sin(\gamma_n x) u_1(t), \qquad (40)$$

where

$$u_{1}(t) = A_{n}e^{\varepsilon_{n}t} + B_{n}e^{\varepsilon_{n}t}$$
$$= \frac{1}{(\varepsilon_{n} - \delta_{n})}[(\lambda VI_{g} - (\delta_{n} - \Omega_{n})I_{f})e^{\varepsilon_{n}t}$$
$$+ (-\lambda VI_{g} + (\varepsilon_{n} - \Omega_{n})I_{f})e^{\delta_{n}t}].$$
(41)

Also, from equation (10), we obtain C(x, t) as

$$C(x, t) = \sum_{n=0}^{\infty} \sin(\gamma_n x) u_2(t), \qquad (42)$$

where

$$u_{2}(t) = E_{n}e^{\varepsilon_{n}t} + F_{n}e^{\delta_{n}t}$$
$$= \frac{1}{(\varepsilon_{n} - \delta_{n})}[(\alpha I_{f} + (\varepsilon_{n} - \Omega_{n})I_{g})e^{\varepsilon_{n}t}$$
$$+ (-\alpha I_{f} + (\Omega_{n} - \delta_{n})I_{g})e^{\delta_{n}t}].$$
(43)

3. Verification at Special Case

In Section 2, the coefficients A_n , B_n , E_n and F_n are obtained in terms of I_f and I_g as

$$A_n = \frac{\lambda V I_g - (\delta_n - \Omega_n) I_f}{\varepsilon_n - \delta_n}, \quad B_n = \frac{-\lambda V I_g + (\varepsilon_n - \Omega_n) I_f}{\varepsilon_n - \delta_n}, \quad (44)$$

$$E_n = \frac{\alpha I_f + (\varepsilon_n - \Omega_n) I_g}{\varepsilon_n - \delta_n}, \quad F_n = \frac{-\alpha I_f + (\Omega_n - \delta_n) I_g}{\varepsilon_n - \delta_n}.$$
 (45)

The main aim of this section is to show that the corresponding coefficients in [9] can be recovered as a special case of the current ones. In [9], the authors considered the ICs (5) as

$$\phi(x, 0) = f(x) = \phi_0, \ C(x, t) = g(x) = h\phi_0, \ 0 < x < L,$$
(46)

where *h* is defined by

$$h = \frac{\beta \nu \sum_{f}}{\lambda} = \frac{\alpha}{\lambda}.$$
(47)

In this case, we can use equation (18) to calculate I_f as

$$I_f = \frac{2}{L} \int_0^L \phi_0 \sin(\gamma_n x) dx = \frac{2\phi_0}{\gamma_n L} [1 - \cos(\gamma_n L)].$$
(48)

From the definition $\gamma_n = (2n+1)\frac{\pi}{L}$, we can find that $\cos(\gamma_n L) = -1$, $\forall n$. Hence, equation (48) reads Developing Ansatz Method for Solving the Neutron Diffusion ... 23

$$I_f = \frac{4\phi_0}{\gamma_n L}.\tag{49}$$

By similar analysis, we can obtain I_g from equation (21) as

$$I_g = \frac{4h\phi_0}{\gamma_n L},\tag{50}$$

i.e., $I_g = hI_f$. Inserting (51) and (52) into (44)-(47), we obtain

$$A_n = \frac{4\phi_0}{\gamma_n L} \left[\frac{\alpha V - \delta_n + \Omega_n}{\varepsilon_n - \delta_n} \right], \quad B_n = \frac{4\phi_0}{\gamma_n L} \left[\frac{-\alpha V + \varepsilon_n - \Omega_n}{\varepsilon_n - \delta_n} \right], \tag{51}$$

$$E_n = \frac{4\phi_0}{\gamma_n L} \left[\frac{\alpha + h(\varepsilon_n - \Omega_n)}{\varepsilon_n - \delta_n} \right], \quad F_n = \frac{4\phi_0}{\gamma_n L} \left[\frac{-\alpha + h(\Omega_n - \delta_n)}{\varepsilon_n - \delta_n} \right].$$
(52)

4. Conclusion

In this paper, a general ansatz method was developed to determine the analytic solution of the neutron flux system under general initial conditions. The neutron flux and the delayed neutron concentration were obtained in explicit forms. The proposed approach is straightforward and also simpler in contrast to other methods in the literature. Moreover, the results in the literature were recovered as a special case of the current ones. Furthermore, the developed approach may deserve a possible extension in future to include complex systems related to neutron diffusion.

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