



THE INVERSE LAPLACE TRANSFORM OF SOME COMPLEX EXPRESSIONS ARISING IN SOLAR ENERGY MODELS

Abdelhalim Ebaid and Amjad A. Alsubaie

Department of Mathematics

Faculty of Science

University of Tabuk

P.O. Box 741, Tabuk 71491

Saudi Arabia

e-mail: aebaid@ut.edu.sa

halimgamil@yahoo.com

441000218@stu.ut.edu.sa

Abstract

Recently, many mathematical models are proposed to describe storage of solar energy. Most of these models are governed by boundary value problems (BVPs). The explicit solutions of such BVPs depend in determining the inverse Laplace transform of complex expressions.

This paper overcomes some of these difficulties arising on account of

Received: May 15, 2023; Accepted: July 1, 2023

2020 Mathematics Subject Classification: 34A25.

Keywords and phrases: Laplace transform, error function, boundary value problem.

How to cite this article: Abdelhalim Ebaid and Amjad A. Alsubaie, The inverse Laplace transform of some complex expressions arising in solar energy models, *Advances in Differential Equations and Control Processes* 30(3) (2023), 297-307.

<http://dx.doi.org/10.17654/0974324323016>

This is an open access article under the CC BY license (<http://creativecommons.org/licenses/by/4.0/>).

Published Online: September 18, 2023

this. The results can be invested to construct the analytic solutions of solar energy models and also models of other fields in engineering sciences.

1. Introduction

In the past decade, many scientific models have been proposed to extract solar energy by means of nanofluids. These models play a vital role in the field of renewable energy which helps in reducing the harmful effects of climate change. This field attracted considerable interest of many researchers [1-7]. The Laplace transform (LT) is a well-known approach to solve initial/boundary value problems (IVPs/BVPs). The LT is of wide application to exactly solve several scientific problems in various fields of science, biology, and engineering [8-16]. It can be seen in these references that the obtained exact/analytic solutions depend mainly on obtaining the inverse LT of complex expressions. So, the objective of this paper is to provide the inversion of some complex expressions which can be invested in future to explore several scientific models.

2. Analysis

This section is devoted to introduce some useful integrals. These integrals are essential to find the inverse LT of complex expressions.

Definition 1. The *complementary error function* (*erfc*) is defined as [17]:

$$\operatorname{erfc}(x) = \frac{2}{\sqrt{\pi}} \int_x^{\infty} e^{-u^2} du. \quad (1)$$

Lemma 1. The integral $I_1 = \int_0^t u^{-3/2} e^{-\left(bu + \frac{a^2}{4u}\right)} du$ is given by

$$I_1 = \frac{\sqrt{\pi}}{a} \left[e^{-a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} - \sqrt{bt}\right) + e^{a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} + \sqrt{bt}\right) \right], \quad a > 0. \quad (2)$$

Proof. Let

$$I_1 = \int_0^t u^{-3/2} e^{-\left(bu + \frac{a^2}{4u}\right)} du. \quad (3)$$

Suppose that

$$z_1 = \sqrt{bu} + \frac{a}{2\sqrt{u}}, \quad z_2 = \sqrt{bu} - \frac{a}{2\sqrt{u}}. \quad (4)$$

Then

$$I_1 = \frac{2}{a} \int_0^t \left[\left(\frac{\sqrt{b}}{2} u^{-1/2} + \frac{a}{4} u^{-3/2} \right) e^{-z_2^2 - a\sqrt{b}} - \left(\frac{\sqrt{b}}{2} u^{-1/2} - \frac{a}{4} u^{-3/2} \right) e^{-z_1^2 + a\sqrt{b}} \right] du, \quad (5)$$

and hence

$$I_1 = \frac{2}{a} \left[\int_{[z_2]_{u=0}}^{[z_2]_{u=t}} e^{-z_2^2 - a\sqrt{b}} dz_2 - \int_{[z_1]_{u=0}}^{[z_1]_{u=t}} e^{-z_1^2 + a\sqrt{b}} dz_1 \right], \quad (6)$$

which is

$$I_1 = \frac{2}{a} \left[e^{-a\sqrt{b}} \int_{-\infty}^{\sqrt{bt} - \frac{a}{2\sqrt{t}}} e^{-z_2^2} dz_2 - e^{a\sqrt{b}} \int_{\infty}^{\sqrt{bt} + \frac{a}{2\sqrt{t}}} e^{-z_1^2} dz_1 \right], \quad a > 0. \quad (7)$$

So

$$I_1 = \frac{2}{a} \left[-e^{-a\sqrt{b}} \int_{\infty}^{\frac{a}{2\sqrt{t}} - \sqrt{bt}} e^{-z_3^2} dz_3 + e^{a\sqrt{b}} \int_{\frac{a}{2\sqrt{t}} + \sqrt{bt}}^{\infty} e^{-z_1^2} dz_1 \right], \quad z_3 = -z_2, \quad (8)$$

i.e.,

$$I_1 = \frac{2}{a} \left[e^{-a\sqrt{b}} \int_{\frac{a}{2\sqrt{t}} - \sqrt{bt}}^{\infty} e^{-z^2} dz_3 + e^{a\sqrt{b}} \int_{\frac{a}{2\sqrt{t}} + \sqrt{bt}}^{\infty} e^{-z_1^2} dz_1 \right]. \quad (9)$$

Thus

$$I_1 = \frac{\sqrt{\pi}}{a} \left[e^{-a\sqrt{b}} \operatorname{erfc} \left(\frac{a}{2\sqrt{t}} - \sqrt{bt} \right) + e^{a\sqrt{b}} \operatorname{erfc} \left(\frac{a}{2\sqrt{t}} + \sqrt{bt} \right) \right], \quad (10)$$

which completes the proof. \square

Lemma 2. For $c > 0$,

$$\begin{aligned} \int_0^t e^{-cz} \operatorname{erfc} \left(\frac{d}{2\sqrt{z}} \right) dz &= -\frac{e^{-ct}}{c} \operatorname{erfc} \left(\frac{d}{2\sqrt{t}} \right) + \frac{1}{2c} \left[e^{-d\sqrt{c}} \operatorname{erfc} \left(\frac{d}{2\sqrt{t}} - \sqrt{ct} \right) \right. \\ &\quad \left. + e^{d\sqrt{c}} \operatorname{erfc} \left(\frac{d}{2\sqrt{t}} + \sqrt{ct} \right) \right]. \end{aligned} \quad (11)$$

Proof. The given integral can be decomposed as

$$\begin{aligned} &\int_0^t e^{-cz} \operatorname{erfc} \left(\frac{d}{2\sqrt{z}} \right) dz \\ &= -\frac{1}{c} \left[e^{-ct} \operatorname{erfc} \left(\frac{d}{2\sqrt{t}} \right) - \frac{d}{2\sqrt{\pi}} \int_0^t z^{-3/2} e^{-\left(cz + \frac{d^2}{4z} \right)} dz \right]. \end{aligned} \quad (12)$$

The integral in the right hand side, in equation (12), can be given from Lemma 1 as

$$\begin{aligned} &\int_0^t z^{-3/2} e^{-\left(cz + \frac{d^2}{4z} \right)} dz \\ &= \frac{\sqrt{\pi}}{d} \left[e^{-d\sqrt{c}} \operatorname{erfc} \left(\frac{d}{2\sqrt{t}} - \sqrt{ct} \right) + e^{d\sqrt{c}} \operatorname{erfc} \left(\frac{d}{2\sqrt{t}} + \sqrt{ct} \right) \right]. \end{aligned} \quad (13)$$

Employing (13) into (12), we have

$$\int_0^t e^{-cz} \operatorname{erfc}\left(\frac{d}{2\sqrt{z}}\right) dz = -\frac{e^{-ct}}{c} \operatorname{erfc}\left(\frac{d}{2\sqrt{t}}\right) + \frac{1}{2c} \left[e^{-d\sqrt{c}} \operatorname{erfc}\left(\frac{d}{2\sqrt{t}} - \sqrt{ct}\right) + e^{d\sqrt{c}} \operatorname{erfc}\left(\frac{d}{2\sqrt{t}} + \sqrt{ct}\right) \right], \quad (14)$$

which is the desired result. \square

Lemma 3. The integral $I_2 = \int_0^t u^{-1/2} e^{-\left(bu + \frac{a^2}{4u}\right)} du$ is given by

$$\frac{\sqrt{\pi}}{2\sqrt{b}} \left[e^{a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} - \sqrt{bt}\right) - e^{-a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} + \sqrt{bt}\right) \right]. \quad (15)$$

Proof. The given integral can be put in the form:

$$I_2 = \frac{1}{\sqrt{b}} \int_0^t u^{-1/2} \left[\left(\frac{\sqrt{b}}{2} u^{-1/2} - \frac{a}{4} u^{-3/2} \right) e^{-z_1^2 - a\sqrt{b}} + \left(\frac{\sqrt{b}}{2} u^{-1/2} + \frac{a}{4} u^{-3/2} \right) e^{-z_2^2 + a\sqrt{b}} \right] du, \quad (16)$$

where z_1 and z_2 are defined by equations (4). Thus

$$I_2 = \frac{1}{\sqrt{b}} \left[\int_{[z_1]_{u=0}}^{[z_1]_{u=t}} e^{-z_1^2 - a\sqrt{b}} dz_1 + \int_{[z_2]_{u=0}}^{[z_2]_{u=t}} e^{-z_2^2 + a\sqrt{b}} dz_2 \right], \quad (17)$$

i.e.,

$$I_2 = \frac{1}{\sqrt{b}} \left[e^{-a\sqrt{b}} \int_{\infty}^{\sqrt{bt} + \frac{a}{2\sqrt{t}}} e^{-z_1^2} dz_1 + e^{a\sqrt{b}} \int_{-\infty}^{\sqrt{bt} - \frac{a}{2\sqrt{t}}} e^{-z_2^2} dz_2 \right], \quad a > 0. \quad (18)$$

Consequently,

$$I_2 = \frac{1}{\sqrt{b}} \left[-e^{-a\sqrt{b}} \int_{\sqrt{bt} + \frac{a}{2\sqrt{t}}}^{\infty} e^{-z_1^2} dz_1 - e^{a\sqrt{b}} \int_{\sqrt{bt} - \frac{a}{2\sqrt{t}}}^{-\infty} e^{-z_2^2} dz_2 \right] \quad (19)$$

or

$$I_2 = \frac{1}{\sqrt{b}} \left[e^{a\sqrt{b}} \int_{\frac{a}{2\sqrt{t}} - \sqrt{bt}}^{\infty} e^{-z_3^2} dz_3 - e^{-a\sqrt{b}} \int_{\frac{a}{2\sqrt{t}} + \sqrt{bt}}^{\infty} e^{-z_1^2} dz_1 \right], \quad z_3 = -z_2, \quad (20)$$

and this gives

$$\begin{aligned} I_2 &= \int_0^{\tau} u^{-1/2} e^{-\left(bu + \frac{a^2}{4u}\right)} du \\ &= \frac{\sqrt{\pi}}{2\sqrt{b}} \left[e^{a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} - \sqrt{bt}\right) - e^{-a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} + \sqrt{bt}\right) \right]. \end{aligned} \quad (21)$$

□

3. Inversion of Complex Expressions

In this section, we obtain inversion of some complex expressions. The current results are of great importance to determine the exact solutions for various IVPs/IVPs that arise in solar energy models.

Lemma 4. For $a \in \mathbb{R}$ (see [17]),

$$\mathcal{L}^{-1}\{e^{-a\sqrt{s+b}}\} = \frac{at^{-3/2}}{2\sqrt{\pi}} e^{-\left(bt + \frac{a^2}{4t}\right)}. \quad (22)$$

Theorem 1. For $a > 0$,

$$\mathcal{L}^{-1}\left\{\frac{e^{-a\sqrt{s+b}}}{s}\right\} = \frac{1}{2} \left[e^{-a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} - \sqrt{bt}\right) + e^{a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} + \sqrt{bt}\right) \right]. \quad (23)$$

Proof. We have

$$\mathcal{L}^{-1}\left\{\frac{e^{-a\sqrt{s+b}}}{s}\right\} = \mathcal{L}^{-1}\left\{\frac{1}{s}\right\} * \mathcal{L}^{-1}\{e^{-a\sqrt{s+b}}\} = \frac{a}{2\sqrt{\pi}} \int_0^t u^{-3/2} e^{-\left(bu + \frac{a^2}{4u}\right)} du, \quad (24)$$

where (*) stands for convolution operation defined by

$$f(t) * g(t) = \int_0^t f(u)g(t-u)du = \int_0^t f(t-u)g(u)du. \quad (25)$$

The integral in equation (24) is already obtained by Lemma 1, hence

$$\mathcal{L}^{-1}\left\{\frac{e^{-a\sqrt{s+b}}}{s}\right\} = \frac{1}{2}\left[e^{-a\sqrt{b}}\operatorname{erfc}\left(\frac{a}{2\sqrt{t}} - \sqrt{bt}\right) + e^{a\sqrt{b}}\operatorname{erfc}\left(\frac{a}{2\sqrt{t}} + \sqrt{bt}\right)\right], \quad (26)$$

and this completes the proof. As a special case, when $b = 0$, the inversion (26) becomes

$$\mathcal{L}^{-1}\left\{\frac{e^{-a\sqrt{s}}}{s}\right\} = \operatorname{erfc}\left(\frac{a}{2\sqrt{t}}\right), \quad (27)$$

which coincides with [17]. □

Lemma 5. *The inversion of $\frac{e^{-d\sqrt{s}}}{s(s-c)}$ is*

$$\begin{aligned} \mathcal{L}^{-1}\left\{\frac{e^{-d\sqrt{s}}}{s(s-c)}\right\} &= -\frac{1}{c}\operatorname{erfc}\left(\frac{d}{2\sqrt{t}}\right) + \frac{1}{2c}\left[e^{ct-d\sqrt{c}}\operatorname{erfc}\left(\frac{d}{2\sqrt{t}} - \sqrt{ct}\right) \right. \\ &\quad \left. + e^{ct+d\sqrt{c}}\operatorname{erfc}\left(\frac{d}{2\sqrt{t}} + \sqrt{ct}\right)\right]. \end{aligned} \quad (28)$$

Proof. The convolution theorem gives

$$\mathcal{L}^{-1}\left\{\frac{e^{-d\sqrt{s}}}{s(s-c)}\right\} = \mathcal{L}^{-1}\left\{\frac{1}{s-c}\right\} * \mathcal{L}^{-1}\left\{\frac{e^{-d\sqrt{s}}}{s}\right\} \quad (29)$$

$$= e^{ct} * \operatorname{erfc}\left(\frac{d}{2\sqrt{t}}\right) = \int_0^t e^{c(t-u)}\operatorname{erfc}\left(\frac{d}{2\sqrt{u}}\right)du \quad (30)$$

$$= e^{ct} \int_0^t e^{-cu}\operatorname{erfc}\left(\frac{d}{2\sqrt{u}}\right)du. \quad (31)$$

Using the result in Lemma 2, we have

$$\int_0^\tau e^{-cu} \operatorname{erfc}\left(\frac{d}{2\sqrt{u}}\right) du = -\frac{e^{-ct}}{c} \operatorname{erfc}\left(\frac{d}{2\sqrt{t}}\right) + \frac{1}{2c} \left[e^{-d\sqrt{c}} \operatorname{erfc}\left(\frac{d}{2\sqrt{t}} - \sqrt{ct}\right) + e^{d\sqrt{c}} \operatorname{erfc}\left(\frac{d}{2\sqrt{t}} + \sqrt{ct}\right) \right]. \quad (32)$$

Employing (32) in (31) leads to the result of this lemma. \square

Lemma 6.

$$\mathcal{L}^{-1}\left\{\frac{e^{-a\sqrt{s+b}}}{s^2}\right\} = \left(\frac{t}{2} e^{-a\sqrt{b}} - \frac{a}{4\sqrt{b}} e^{a\sqrt{b}}\right) \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} - \sqrt{bt}\right) + \left(\frac{\tau}{2} e^{a\sqrt{b}} - \frac{a}{4\sqrt{b}} e^{-a\sqrt{b}}\right) \operatorname{erfc}\left(\frac{a}{2\sqrt{\tau}} + \sqrt{b\tau}\right). \quad (33)$$

Proof. Applying the convolution theorem, we have

$$\mathcal{L}^{-1}\left\{\frac{e^{-a\sqrt{s+b}}}{s^2}\right\} = \mathcal{L}^{-1}\left\{\frac{1}{s^2}\right\} * \mathcal{L}^{-1}\{e^{-a\sqrt{s+b}}\}, \quad (34)$$

$$\begin{aligned} \mathcal{L}^{-1}\left\{\frac{e^{-a\sqrt{s+b}}}{s^2}\right\} &= \frac{a}{2\sqrt{\pi}} \left(t * t^{-3/2} e^{-\left(bt + \frac{a^2}{4t}\right)} \right) \\ &= \frac{a}{2\sqrt{\pi}} \int_0^t (t-u) u^{-3/2} e^{-\left(bu + \frac{a^2}{4u}\right)} du \\ &= \frac{a}{2\sqrt{\pi}} (tI_1 - I_2), \end{aligned} \quad (35)$$

where I_1 and I_2 were already obtained by Lemma 1 and Lemma 3, respectively, i.e.,

$$\begin{aligned}
 I_1 &= \int_0^t u^{-3/2} e^{-\left(bu + \frac{a^2}{4u}\right)} du \\
 &= \frac{\sqrt{\pi}}{a} \left[e^{-a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{u}} - \sqrt{bt}\right) + e^{a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} + \sqrt{bt}\right) \right], \quad (36)
 \end{aligned}$$

while

$$\begin{aligned}
 I_2 &= \int_0^t u^{-1/2} e^{-\left(bu + \frac{a^2}{4u}\right)} du \\
 &= \frac{\sqrt{\pi}}{2\sqrt{b}} \left[e^{a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} - \sqrt{bt}\right) - e^{-a\sqrt{b}} \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} + \sqrt{bt}\right) \right]. \quad (37)
 \end{aligned}$$

Employing I_1 and I_2 into equation (35) and simplifying, we have

$$\begin{aligned}
 \mathcal{L}^{-1} \left\{ \frac{e^{-a\sqrt{s+b}}}{s^2} \right\} &= \left(\frac{t}{2} e^{-a\sqrt{b}} - \frac{a}{4\sqrt{b}} e^{a\sqrt{b}} \right) \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} - \sqrt{bt}\right) \\
 &\quad + \left(\frac{t}{2} e^{a\sqrt{b}} - \frac{a}{4\sqrt{b}} e^{-a\sqrt{b}} \right) \operatorname{erfc}\left(\frac{a}{2\sqrt{t}} + \sqrt{bt}\right). \quad (38)
 \end{aligned}$$

□

4. Conclusions

The explicit solutions of many mathematical models describing storage of solar energy depend on evaluating the inverse Laplace transform of complex expressions. In this paper, some of the difficulties arising due to it are solved. Accordingly, the obtained results are useful in establishing explicit solutions of several scientific models including solar energy and also other areas of engineering sciences.

References

- [1] Y. R. Sekhar, K. V. Sharma and S. Kamal, Nanofluid heat transfer under mixed convection flow in a tube for solar thermal energy applications, *Environ. Sci. Pollut. Res.* 23 (2016), 9411-9417. <https://doi.org/10.1007/s11356-015-5715-9>.
- [2] N. A. Sheikh, F. Ali, I. Khan and M. Gohar, A theoretical study on the performance of a solar collector using CeO_2 and Al_2O_3 water based nanofluids with inclined plate: Atangana-Baleanu fractional model, *Chaos Solitons Fractals* 115 (2018), 135-142.
- [3] R. Mehmood, Rabil Tabassum, S. Kuharat, O. Anwar Bég and M. Babaie, Thermal slip in oblique radiative nano-polymer gel transport with temperature-dependent viscosity: solar collector nanomaterial coating manufacturing simulation, *Arabian Journal for Science and Engineering* 44 (2019), 1525-1541. <https://doi.org/10.1007/s13369-018-3599-y>.
- [4] A. M. Norouzi, M. Siavashi and M. Khaliji Oskouei, Efficiency enhancement of the parabolic trough solar collector using the rotating absorber tube and nanoparticles, *Renew. Energy* 145 (2020), 569-584. <http://dx.doi.org/10.1016/j.renene.2019.06.027>.
- [5] S. E. Ghasemi and M. Hatami, Solar radiation effects on MHD stagnation point flow and heat transfer of a nanofluid over a stretching sheet, *Case Studies in Thermal Engineering* 25 (2021), 100898. <https://doi.org/10.1016/j.csite.2021.100898>.
- [6] E. A. C. Panduro, F. Finotti, G. Largiller and K. Y. Lervåg, A review of the use of nanofluids as heat-transfer fluids in parabolic-trough collectors, *Applied Thermal Engineering* 211 (2022), 118346.
- [7] A. F. Aljohani, A. Ebaid, E. H. Aly, I. Pop, A. O. M. Abubaker and D. J. Alanazi, Explicit solution of a generalized mathematical model for the solar collector/photovoltaic applications using nanoparticles, *Alexandria Engineering Journal* 67 (2023), 447-459. <https://doi.org/10.1016/j.aej.2022.12.044>.
- [8] A. Ebaid and M. Al Sharif, Application of Laplace transform for the exact effect of a magnetic field on heat transfer of carbon-nanotubes suspended nanofluids, *Z. Nature. A* 70(6) (2015), 471-475.
- [9] A. Ebaid, A. M. Wazwaz, E. Alali and B. Masaedeh, Hypergeometric series solution to a class of second-order boundary value problems via Laplace transform with applications to nanofluids, *Commun. Theor. Phys.* 67 (2017), 231.

- [10] H. Saleh, E. Alali and A. Ebaid, Medical applications for the flow of carbon-nanotubes suspended nanofluids in the presence of convective condition using Laplace transform, *J. Assoc. Arab Univ. Basic Appl. Sci.* 24 (2017), 206-212.
- [11] A. Ebaid, E. Alali and H. Saleh, The exact solution of a class of boundary value problems with polynomial coefficients and its applications on nanofluids, *J. Assoc. Arab Univ. Basic Appl. Sci.* 24 (2017), 156-159.
- [12] S. M. Khaled, The exact effects of radiation and joule heating on magnetohydrodynamic Marangoni convection over a flat surface, *Therm. Sci.* 22 (2018), 63-72.
- [13] H. S. Ali, E. Alali, A. Ebaid and F. M. Alharbi, Analytic solution of a class of singular second-order boundary value problems with applications, *Mathematics* 7 (2019), 172. <https://doi.org/10.3390/math7020172>.
- [14] A. Ebaid, W. Alharbi, M. D. Aljoufi and E. R. El-Zahar, The exact solution of the falling body problem in three-dimensions: comparative study, *Mathematics* 8(10) (2020), 1726. <https://doi.org/10.3390/math8101726>.
- [15] S. M. Khaled, A. Ebaid and F. Al Mutairi, The exact endoscopic effect on the peristaltic flow of a nanofluid, *J. Appl. Math.* 2014 (2014), 11, Article ID 367526. <http://dx.doi.org/10.1155/2014/367526>.
- [16] A. F. Aljohani, A. Ebaid, E. A. Algehyne, Y. M. Mahrous, P. Agarwal, M. Areshi and H. K. Al-Jeaid, On solving the chlorine transport model via Laplace transform, *Scientific Reports* 12 (2022), 12154. <https://doi.org/10.1038/s41598-022-14655-3>.
- [17] M. R. Spiegel, *Laplace Transforms*, McGraw-Hill, Inc., 1965.